

# Oxygen transport and release of adenosine triphosphate in micro-channels and arterioles in the human microcirculation

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## Abstract

The governing nonlinear equations for oxygen transport and ATP concentration in a microfluidic channel and tube are solved in a novel way with the aid of Maple and COMSOL Multiphysics simulation software. Considering a model which assumes a larger plasma velocity near the wall of a channel/tube in comparison to RBC velocity near the centerline we obtain results showing clearly that there is a significant decrease in oxygen tension in the vicinity of an oxygen permeable membrane placed midway on the upper channel/tube wall and to the right side of it in the downstream field. The purpose of the membrane is to cause a rapid change in oxygen saturation as RBC's flow through channel/tube. To the right of the membrane downstream the greatest amount of ATP is released. It is shown that for smaller arterioles there is an optimal size for which greater amounts of ATP can be released. Finally the corresponding time-dependent oxygen transport problem for plug flow in a channel, which has not been simulated in previous models in the literature, is analyzed and different starting times are shown for ATP release at different points in the channel. The FE modelling is very accurate: The time evolution problem is modelled and solved in its entirety with exact parameters used in the literature for blood flow and oxygen transport in the microcirculation. A comparison is made between the steady state and time dependent solutions in order to validate the results. The implications of the time dependent model for biological systems such as the human microcirculation requires exact information on release of energy as ATP is released from blood cells and the present work is important in providing this information. In particular, the time varying problem for ATP concentration and its temporal dynamics, as calculated in this paper is crucial in determining the ADP/ATP concentration pair which is an oscillating pair in the biochemical oscillations that occur in glycolysis in anaerobic yeast cultures.

# 1.Introduction

ATP is made up of Adenine, Ribose and Phosphate groups. It is the molecule of most organisms to transfer energy. Adenine is one of the bases found in DNA and RNA. Ribose is a sugar molecule found in RNA. There exist high energy bonds between the phosphates. Breaking the first high energy bond in ATP releases energy and a free phosphate. This energy is used to power cellular activities. What is left with 2 phosphates and Adenosine is called Adenine diphosphate or ADP. Energy from the breakdown of food is used to change the ADP to an ATP molecule. In cells, mitochondria supply the major portion of ATP consumed in the cytosol. ADP-ATP translocation in mitochondria is studied by Klingenberg [1]. Mitchell [2] proposed that electron transport and ATP synthesis are related by a proton gradient across the inner mitochondrial membrane. A mechanism for microvascular blood flow regulation that is under current investigation depends on the oxygen saturation-dependent release by red blood cells (RBCs) of the signaling molecule ATP [3,4]. Increased RBC ATP release due to decreased O<sub>2</sub> saturation results in a vasodilation signal that conducts upward through the microvascular tree. This signal can originate in arterioles, capillaries or venules. It is arterioles that the present work focuses on. Moschandreou et al. [5] have presented a theoretical model for steady-state radial and longitudinal oxygen transport in arterioles containing plasma and red blood cells and surrounded by living tissue. The governing equations were solved in a cell-rich and cell-free region. Using the transport equations for O<sub>2</sub> transport in [5] together with ATP concentration equations in [6,7], both defined in a micro-channel and then applied to representative arterioles, important information is obtained about ATP release. 3D printed microfluidic devices have been recently used to collect ATP while simultaneously measuring the release stimulus which is the reduced oxygen concentration [8,9]. The potential for the technology to significantly impact the field of microfluidics is very good. [10,11] In [11], well inserts in a microfluidic device were filled with 200  $\mu\text{L}$  of Hanks' balanced salt solution (HBSS) and stable RBC flow ( $6\mu\text{L min}^{-1}$ ) was established with no air bubbles. Efforts were made to minimize evaporation from wells and 7% RBCs flowed for 20 minutes. RBC-derived ATP was measured when a luciferin luciferase mixture was applied and resulting chemiluminescence was determined.

In the present paper a novel solution method is used to determine the oxygen tension and concentration of ATP in a microfluidic channel with dimensions, inlet and boundary conditions given by Sove et al.[6], then solve the governing equations of oxygen and ATP transport in tubes of different diameters and finally present the corresponding time-dependent solution in a channel with plug flow. Three objectives of the present work is to provide a compact and accessible method of solution in Maple[12], secondly to include a realistic model of blood flow in arterioles(tubes) calculating the oxygen transport and concentration of ATP produced due to a change in oxygen saturation and finally solve the corresponding time dependent problem using COMSOL Multiphysics [13] which will predict the onset times of ATP in the channel.

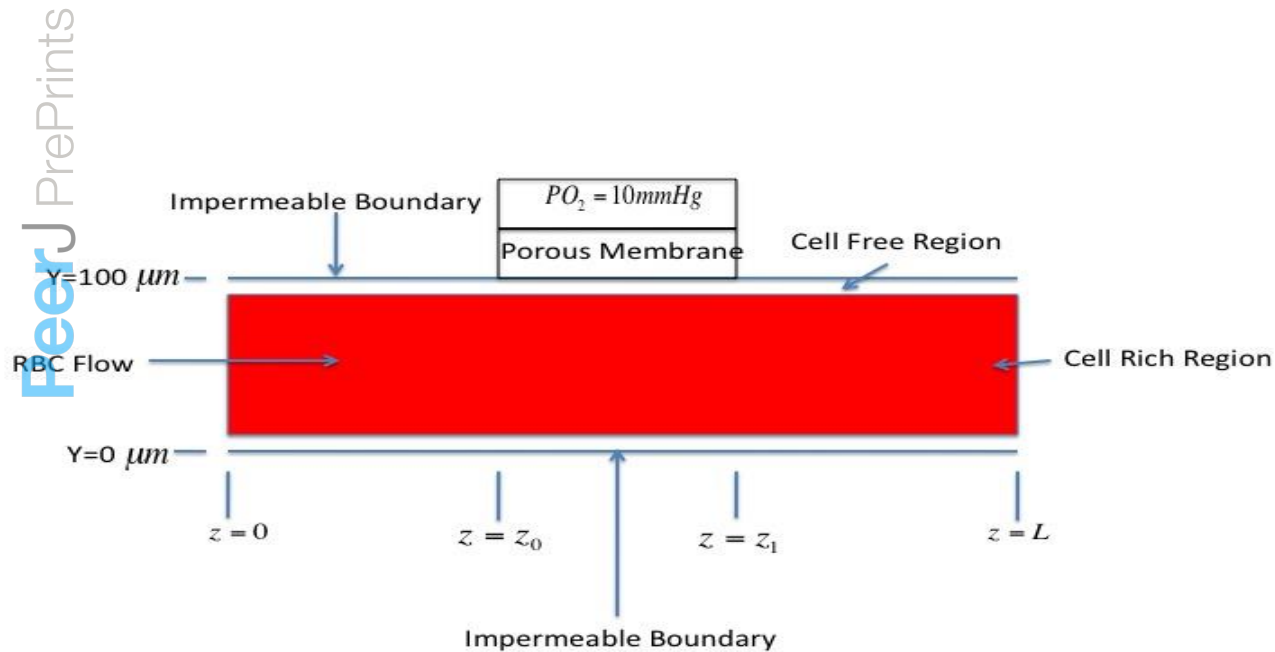
## 2. Materials and Methods

### 2.1 Channel Flow and Oxygen Transport

Hellums and co-workers [14,15,16] have developed a model for intraluminal oxygen transport in arteriole size vessels.

We adopt this model for square channel flow where Equation (1) is used below to model oxygen transport in a micro-fluidic device.

The geometry of the present mathematical model is shown in Figure 1.



**Figure 1**

Geometry of mathematical model used for channel flow

$$\left( v_p (1 - H_T) + v_{RBC} H_T \frac{K_{RBC}}{K_p} \left( 1 + \frac{[Hb_T]}{K_{RBC}} \frac{dSO_2}{dPO_2} \right) \right) \frac{\partial PO_2}{\partial z} = D_p \frac{\partial^2 PO_2}{\partial y^2} \quad (1)$$

There is a core region of blood flow with RBCs and plasma and a cell-free region with only plasma flowing,  $v_p$  is the blood plasma velocity and  $v_{RBC}$  is the velocity of RBCs with plasma in the cell rich region in Figure 1. The velocity of the RBCs is lower due to the slip between plasma and RBCs. [17] The distribution of RBC's is such that hematocrit is higher at the center of the channel and lower near the wall. [ 5 ] .Dynamic hematocrit reduction is known as the Fahraeus effect. [18 ] In blood vessels with diameters less than 500 microns both the hematocrit decreases with decreasing vessel diameter. There is a reduction of red cell transit time through a given tube segment known as tube hematocrit  $H_T$ . The discharge hematocrit, is the hematocrit of blood entering or leaving the tube segment.

The term  $\frac{dSO_2}{dPO_2}$  is the slope of the oxyhemoglobin dissociation curve and is a nonlinear function of  $PO_2$ . (Nair et al. [15,16 ]) The dissociation curve is approximated by the Hill equation in Eq. (1a) , where N is an empirical constant and  $P_{50}$  is the oxygen tension that yields 50 % oxygen saturation.  $[Hb_T]$  is the total heme concentration which is equal to four times the hemoglobin concentration due to fact that there are four heme groups on each hemoglobin molecule.  $D_p$  is the oxygen diffusion coefficient in plasma and has units of  $\mu m^2 / s$ , and  $K_{RBC}$  and  $K_p$  are the solubilities of  $O_2$  in the RBCs and plasma, respectively, and have units of M/mmHg. Differentiating the Hill equation with respect to  $PO_2$  gives Eq. (1b).

$$SO_2(y, z) = \frac{PO_2^N(y, z)}{PO_2^N(y, z) + P_{50}^N} \quad (1a)$$

$$\frac{dSO_2}{dPO_2} = \frac{NP_{50}^N P^{N-1}}{(P_{50}^N + P^N)^2} \quad (1b)$$

Generally blood flow velocity varies from zero at the wall of the channel to maximum at center. In this paper we have assumed two types of flow, one plug flow and the other Poiseuille flow.

In [19] a quoted value of arteriole flow velocity is in the range of 1 to 1.275 mm/s (1000 to 1275 microns/s). For discharge hematocrits of about 0.4 we have used a value of 0.7 for the ratio of tube hematocrit to discharge hematocrit.[18,20] These values in [18,20] are assumed for arterioles of diameter 20 down to 10 microns. We have used a flow velocity of 1275 for  $v_p$  and a representative value for discharge hematocrits of 0.4 for channel and tube flow ranging from 100 down to 10 microns radius.

$$\left( 1275 \cdot 0.72 + \frac{1147.5 \cdot 0.28 \cdot 1.47}{1.33} \times \left( 1 + \frac{5350 \cdot 2.7 \cdot 27^{2.7} \cdot PO_2(y, z)^{1.7}}{1.47 \cdot \left( PO_2(y, z)^{2.7} + 27^{2.7} \right)^2} \right) \right) \frac{\partial PO_2}{\partial z} = 2750 \cdot \frac{\partial^2 PO_2}{\partial y^2} \quad (2)$$

## 2.2 Channel Flow and ATP Model

ATP concentration given by  $C_{ATP}$  is governed by the following transport equation.

$$v_p \frac{\partial C_{ATP}}{\partial z} (1 - H_T) = D_{ATP} \frac{\partial^2 C_{ATP}}{\partial y^2} (1 - H_T) + R(H_T) - k_t C_{ATP} (1 - H_T) \quad (3)$$

where  $D_{ATP}$  is the diffusion coefficient of ATP in blood plasma with units of  $\mu m^2 / s$  and  $k_t$  is a concentration rate constant with units of  $s^{-1}$ .

Equations (1) and (3) are coupled partial differential equations through the nonlinear function  $R(y, z)$  [6],[21],

$$R(y, z) = -14 \cdot SO_2(y, z) + 14 \quad (4)$$

where,

$$SO_2(y, z) = \frac{PO_2^{2.7}(y, z)}{PO_2^{2.7}(y, z) + 3600^{2.7}} \quad (5)$$

and a maximum ATP release rate of  $14 \mu M / s$  .[6]

Equations 1 and 3 are solved using Maple 13 [12].

### 2.3 Channel Inlet and Boundary Conditions for Oxygen Transport and ATP Concentration

The inlet conditions for both oxygen tension and ATP concentration at the entrance to the channel where the oxygen porous membrane is located at  $z=7000$  microns are given by Equation 6:

$$PO_2(y,7000) = P_0 = 150 \text{ mmHg}$$

$$C_{ATP}(y,7000) = C_0 = 0 \mu M \tag{6}$$

The boundary conditions for both oxygen tension and ATP concentration in the channel are given by Equation 7:

$$\frac{\partial PO_2}{\partial y}(h, z) = 0 \quad 0 \leq z < z_0, \quad z_1 < z \leq L$$

$$\frac{\partial PO_2}{\partial y}(0, z) = 0 \quad 0 \leq z \leq L$$

$$D_p K_p \frac{\partial PO_2}{\partial y}(h, z) = D_m K_m \frac{PO_{2(PM)} - PO_2(h, z)}{\tau} \quad z_0 \leq z \leq z_1 \quad (7)$$

$$\frac{\partial C_{ATP}}{\partial y}(h, z) = 0 \quad z_0 \leq z \leq z_1$$

$$\frac{\partial C_{ATP}}{\partial y}(0, z) = 0 \quad z_0 \leq z \leq z_1$$

### TABLE 1 Model Parameters

Physical parameters used in oxygen transport model.

Blood plasma velocity	$v_p (\mu m / s)$	1275[19]
Red Blood Cell velocity	$v_{RBC} \mu m / s$	1147.5 [19]
Slip coefficient	Slp	0.1 [17]
Discharge Hematocrit	$H_D$	0.4 [18,20]
Tube Hematocrit	$H_T$	0.28 [18,20]
$O_2$ Diffusivity in plasma	$D_p (\mu m^2 / s)$	2750[5],[6]
$O_2$ Diffusivity in $O_2$ Permeable Membrane	$D_m (\mu m^2 / s)$	160000[6]
Plasma Layer Thickness	$T(\mu m)$	1[6]

$O_2$ Permeable Membrane $PO_2$	$PO_{2(PM)}$ (mmHg)	10[6]
Axial position for start of $O_2$ Permeable Membrane along channel/arteriole	$z_0$ ( $\mu m$ )	7000 [6]
Axial position for end of $O_2$ Permeable Membrane along channel/arteriole	$z_1$ ( $\mu m$ )	7700[6]
Inlet $PO_2$	$P_0$ (mmHg)	150 [6]
Microfluidic Device/Arteriole length	$L$ ( $\mu m$ )	14700
$O_2$ Porous Membrane Thickness	$\tau$ ( $\mu m$ )	100 [6]
Hill Coefficient	$N$	2.7[5]
$O_2$ Solubility in Porous Membrane	$K_m$ ( $\frac{\mu M}{mmHg}$ )	17.959[6]
$O_2$ Solubility in Plasma	$K_p$ ( $\frac{\mu M}{mmHg}$ )	1.33[5],[6]
$O_2$ Solubility in RBCs	$K_{RBC}$ ( $\frac{\mu M}{mmHg}$ )	1.47 [5],[6]
Reaction Rate	$K_t$ $s^{-1}$	0.1 [6]
Total Heme Concentration	$[Hb_T]$ ( $\mu M$ )	5350 [5],[6]



ATP Diffusivity in plasma	$D_{ATP} (\mu m^2 / s)$	475 [6]
Partial Pressure at 50% Saturation	$P_{50} (mmHg)$	27 [5],[6]

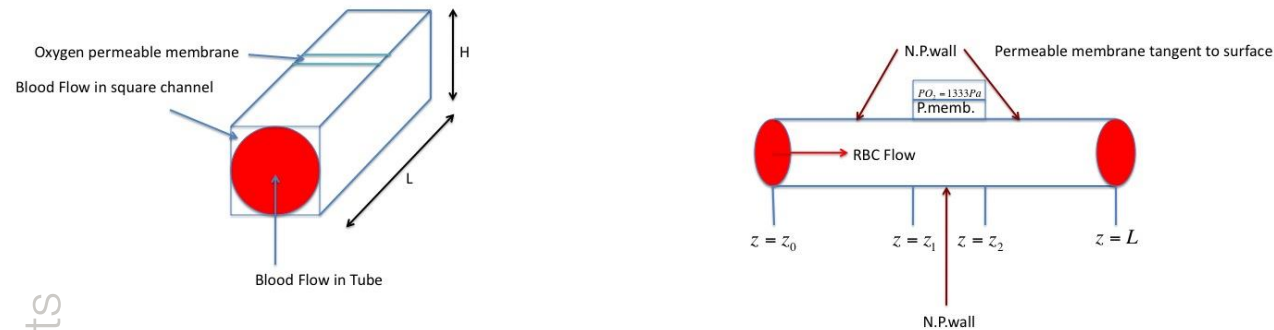
## 2.4 Tube Flow and Oxygen Transport and ATP Concentration

In arteriolar(tube) flow Equation (8) is used below to model oxygen transport. The model used includes a realistic diffusive oxygen transport term.

Whereas a channel is considered in [6], the inclusion of an extra radial term is used in the present work to effectively model arteriolar oxygen transport.

$$\left( v_p (1 - H_T) + v_{RBC} H_T \frac{K_{RBC}}{K_p} \left( 1 + \frac{[Hb_T] dSO_2}{K_{RBC} dPO_2} \right) \right) \frac{\partial PO_2}{\partial z} = D_p \left( \frac{\partial^2 PO_2}{\partial y^2} + \frac{1}{y} \frac{\partial PO_2}{\partial y} \right) \quad (8)$$

Same inlet and boundary conditions are used here as in channel flow in section 2.3. A tangent surface on the top of tube defines the boundary condition as in channel flow which acts along a line on the top of the surface of the tube as shown in Figure 2.



**Figure 2**

Geometry of Model Used for Tube Flow

ATP concentration in a tube given by  $C_{ATP}$  is governed by the following transport equation.

$$v_p \frac{\partial C_{ATP}}{\partial z} (1 - H_T) = D_{ATP} \left( \frac{\partial^2 C_{ATP}}{\partial y^2} + \frac{1}{y} \frac{\partial C_{ATP}}{\partial y} \right) (1 - H_T) + R(H_T) - k_t C_{ATP} (1 - H_T) \quad (9)$$

where  $D_{ATP}$  is the diffusion coefficient of ATP in blood plasma and  $k_t$  is the concentration rate constant defined previously.

## 2.5 Parabolic Flow in Tubes of Radius R=10, 20 Microns

Convective transport in the cell-free (plasma) and cell-rich (RBCs) regions involves specific assumed velocity profiles.

These are given as parabolic profiles and thus Poiseuille flow is assumed to occur in the tube shown in Figure 2 . The following equations are used for blood flow velocity at the inlet of the tube with h=10 and 20 microns respectively.

$$v_p = 1275 \cdot (h^2 - y^2) \quad (10)$$

$$v_{RBC} = 1147.5 \cdot (h^2 - y^2) \quad (11)$$

$$\left( v_p (1 - H_T) + v_{RBC} H_T \frac{K_{RBC}}{K_p} \left( 1 + \frac{[Hb_T]}{K_{RBC}} \frac{dSO_2}{dPO_2} \right) \right) \frac{\partial PO_2}{\partial z} = D_p \left( \frac{\partial^2 PO_2}{\partial y^2} + \frac{1}{y} \frac{\partial PO_2}{\partial y} \right) \quad (12)$$

The boundary conditions used in the simulations are given by Eqs. (6) and (7) with a note that the radial derivative involves a singular coefficient and is treated separately in the Maple coding. Here the derivative is at y=0 where y denotes the radial coordinate in Eq. (12).

## 2.6 Time Dependent Channel Flow

The governing equations for time dependent erythrocytic flow in a microfluidic channel are given by the following[22,23] :

$$\left( H_T C_{Hb} \frac{\partial SO_2}{\partial PO_2} \right) \frac{\partial PO_2}{\partial t} + \left( v_p (1 - H_T) + v_{RBC} H_T \frac{K_{RBC}}{K_p} \cdot \left( 1 + \frac{C_{Hb}}{1 \times 10^{-4}} \right) \right) \cdot \frac{\partial SO_2}{\partial PO_2} \quad (13)$$

$$\times \frac{\partial PO_2}{\partial z} = \frac{1}{R^2} \cdot j_{wall} + \frac{1}{R^2} \cdot D_p \cdot \frac{\partial^2 PO_2}{\partial y^2} = \frac{1}{R^2} \cdot \frac{\partial PO_2}{\partial y} \Big|_{y=H} + \frac{1}{R^2} \cdot D_p \cdot \frac{\partial^2 PO_2}{\partial y^2}$$

$$\left( 0.6 \cdot 0.5 \times 10^8 \cdot H_T C_{Hb} \right) \cdot \frac{\partial SO_2}{\partial PO_2} \cdot \frac{\partial PO_2}{\partial t} + \left( 1275 \cdot 0.72 + 1147.5 \cdot 0.28 \cdot 1.47 / 1.33 \cdot \left( \frac{0.208}{1 \times 10^{-4}} \right) \right) \frac{\partial SO_2}{\partial PO_2} \quad (14)$$

$$\times \frac{\partial PO_2}{\partial z} = \frac{1}{100^2} \cdot j_{wall} + \frac{1}{100^2} \cdot 2750 \cdot \frac{\partial^2 PO_2}{\partial y^2} = \frac{1}{100^2} \cdot \frac{\partial PO_2}{\partial y} \Big|_{y=H} + \frac{1}{100^2} \cdot 2750 \cdot \frac{\partial^2 PO_2}{\partial y^2}$$

COMSOL Multiphysics [13] is used to handle Eqs. (13) and (14).

Equation (13) and (14) have been written slightly different than Eq. (1) in that we have used oxygen binding capacity, added a flux term as in [22,23] and added a time derivative and coefficient. R is set to 100 microns.

It is assumed that for Hemoglobin(whole blood) that the conversion factor is on average 15 g/dL [24].  $C_{Hb}$  is the  $O_2$  - binding capacity of RBC Hb solution.  $C_{Hb}$  is the product of the RBC hemoglobin concentration and the  $O_2$ -binding capacity of Hb. The oxygen binding capacity of human haemoglobin in [25] has been reported to be in the mean  $1.368 \text{ ml} \cdot \text{g}^{-1}$ .  $C_{Hb}$  is approximately equal to 0.208. This is used in COMSOL's interface for the time derivative coefficient which is in units of seconds. In order to set the mass coefficient to units of seconds, we rewrote the units as ,

$$\frac{\mu\text{m}^2}{\text{s}} \cdot \frac{\text{s}^3}{\text{kg}} = \frac{\mu\text{m}^2}{\text{s}} \cdot \frac{1}{\text{kg/s}} \cdot \text{s}^2$$

and used a flow rate of 7.8 microlitres/min [6]. Converting to kg/s a factor of approximately  $0.5 \times 10^8$  was introduced in mass coefficient in COMSOL. This was set since the units of  $\Gamma$  and  $f$  (Appendix B) were independent of kg/s.(ie cancellation of pascal units). We have multiplied the product of tube hematocrit and binding capacity in the time derivative coefficient by a factor of 0.6 matching the initial time of 25 ms release time reported in [6], occurring downstream in channel.

### 3. Numerical Methods

#### 3.1 Maple Solution Scheme

The default method is used in Maple 13 [12] which uses a second order (in radial or y direction and axial direction z) centered, implicit finite difference scheme to obtain the solution. The number of points in the stencil of finite difference scheme is one greater than the order of each equation. A first order accurate boundary condition is used for a second order accurate method and still produces a solution of second order accuracy.

Derivatives in the initial/boundary conditions are specified in indexed D notation. ( $D_p K_p D_1(PO_2)(20, x) = D_m K_m \frac{10 - PO_2(20, x)}{\tau}$  is the boundary condition along oxygen porous membrane.) The calling sequence in Maple is `pdsolve(PDE, conditions, numeric, options)`. The parameters include the type of PDE which allows for a single z-dependent partial differential equation in one independent variable y, and the set of initial and boundary conditions given by Eq. (6) and (7). A numeric keyword is used to invoke a centered, implicit finite difference scheme and options specifies the spacing of the spatial points on the discrete mesh on which solution computes and defaults to 1/20 th of the spatial range of problem.

#### Case 1

In maple the following commands are used.

```
pds:= pdsolve(sys,IBC,numeric)
```

$$IBC := \{PO_2(y, 7000) = 150, D_1(PO_2)(0, z) = 0, D_1(PO_2)(100, z) = 787 \cdot 0.1 - 787 \cdot 0.01 \cdot PO_2(100, z), \\ C_{ATP}(y, 7000) = 0, D_1(C_{ATP})(100, z) = 0, D_1(C_{ATP})(0, z) = 0\} \quad (15)$$

The boundary conditions and entry conditions are such that the entry condition for blood flow is uniform oxygen tension of 150 mmHg at a length of 7000 microns from inlet of microfluidic device and a flux condition through permeable membrane shown in Figure 1, due to oxygen source at permeable wall at length  $z = z_0 = 7000$  microns, height  $y = 100$  microns. [6] The oxygen flux through membrane was set above in definition of IBC between  $z_0$  and  $z_1$  in Fig(1).

The condition of oxygen impermeability between 0 and  $z_0$  and  $z_1$  and L is set in the IBC definition in Maple for both oxygen tension and ATP concentration. The novelty of present work is that in Maple the command `pdsolve` was used again to solve Equation 1 using the solution in Case 1 above as an entry condition for the region between  $z_1$  and L (Fig 1) where a no flux condition or flux of zero is set in the definition of the boundary condition.

That is we have the following :

#### Case 2.

In Maple the solution obtained from Case1 is written as:

$$U:=\text{subs}(pds:-\text{value}(\text{output}=\text{listprocedure}), PO_2(y, z) ); \quad (16)$$

Proceeding in Maple language, Eq. 1 is defined as in Case 1 by defining the equation by a variable 'PDE'.

The entry and boundary conditions become as follows using the command given by Eq. (16):

$$IBC := \{PO_2(y, 7700) = U(y, 7700), D_1(PO_2)(0, z) = 0, D_1(PO_2)(100, z) = 0$$

It can be observed that the solution obtained from Case 1 is used up to the end of the oxygen permeable wall and set to be an inlet condition in  $z$  to where the non-permeable wall condition starts to the right of the oxygen porous membrane. The advantage to store the solution in the computational field by using (16) is to solve the partial differential equations in two steps in the entire region from  $z=0$  to  $z=L$  as shown in Figure 1.

### 3.2 COMSOL Multiphysics Solution Scheme

The General Form PDE Interface in COMSOL Multiphysics 4.3b [ 13 ] was used in time dependent simulations where the initial/boundary conditions were implemented as shown in Equation (6) and (7) with the addition of an initial condition in time ( $t=0$ ) for  $PO_2$ , equal to 20kPa or equivalently 150 mmHg defined at every point in the channel. The general form is used in COMSOL for nonlinear PDE's and is appropriately set as follows. After piecing together three rectangles in COMSOL geometry interface due to permeable membrane and nonpermeable walls, the general form PDE defined on resulting domain is [13] shown in Appendix B.

Choosing options from PDE interface as ``Zero Flux`` and ``Flux/Source`` as well as ``Initial Values`` all boundary/initial values were specified as in Equation (6) and (7) above and condition at  $t=0$  seconds.

As COMSOL Multiphysics is a finite-element method based software package[13], a physics-controlled mesh with extra fine element size was used in all simulations.

Finally the source term above was set to  $f$  and  $e_a = 0$ .  $\Gamma_x$  and  $\Gamma_y$  have units of N/m,  $f$  has units of Pa, and  $d_a$  has units of seconds. The boundary condition in COMSOL at oxygen permeable membrane is defined as,

$$-\mathbf{n} \cdot \Gamma = g - qu$$

(17)

where  $g$  is a boundary flux/source term and  $q$  is a boundary absorption term.

For the time dependent problem, the boundary condition at permeable membrane is,

$$D_p K_p \frac{\partial PO_2}{\partial y}(h, z, t) = D_m K_m \frac{PO_{2(PM)} - PO_2(h, z, t)}{\tau} \quad z_0 \leq z \leq z_1$$

and thus we take the same boundary condition in the COMSOL time dependent problem as in maple solution,

$$g = 78.7 \times 1333.33$$

$$q = 7.87$$

With these definitions Equations (1,13-14) , boundary conditions (6,7) and initial condition of 20 kPa at  $t=0$  were input to COMSOL interface.

## 4. Results and Discussion

### Steady State and Time-Dependent Oxygen Transport and ATP Release in Micro-Channel and Tube

The computational model of  $O_2$  transport and ATP concentration in a micro-channel and in an arteriole includes a more physiological equation that has been used before, that is, a radial and time-dependent transport equation. For different size arterioles, we calculated how the radially changing  $PO_2$  inside the vessels decreases with increasing distance  $z$  downstream from the vessel inlet. Our calculations show appreciable drops in radial intraluminal  $PO_2$  downstream from where the oxygen permeable boundary condition is set [6]. These results are similar to the results in [6], however in addition include realistic tube-like geometries, differently assumed velocity profiles and time dependent oxygen transport.

For steady-state oxygen transport in a channel of radius 100 microns it is observed that the greatest decrease in oxygen tension is near the oxygen permeable membrane and to the right of it downstream and in Figure 3-A an ATP concentration as large as 0.08 micro-moles is shown near the membrane. For values further downstream, the ATP concentration flattens out across channel as observed in Figure 3-B and converges to approximately 1 micro-mole as distance tends to very large axial values. The oxygen tension decreases to a greater extent at the wall of channel at the end of permeable membrane at 7700 microns than at impermeable end  $y=0$ . For values greater than this downstream, the oxygen tension decreases to the same extent near the wall than at 7700 microns. In addition to this the profiles tend to flatten out in  $y$ -direction near axial distance 11700 microns and greater. In Figure 3-C the value obtained is approximately 13345 Pa (100 mmHg) in the uniform  $y$ -direction and at the end of channel. The inlet value is in accordance to Sove's work [6] with a value of 150 mmHg and our results for the steady state problem in the channel is slightly more than Sove's value of 95 mmHg uniform across the height of channel and downstream at the end of the channel at  $z=L$ . This is different since we have assumed a different velocity profile in comparison to [6]. Using the velocity from Sove et al. [6] [see Appendix], we obtained in Maple approximately 94 mm Hg uniform across height of channel which is closer to their results, however the same simulation was performed in COMSOL and we obtained 94 mm Hg. It is note worthy that this agreement for steady state results between Maple and COMSOL simulations is important in order to compare results reliably to the time dependent problem below using COMSOL, especially since Maple's `pdsolve` command is not capable of handling problems with two dimensions and time.

It is concluded in Figure 3 that there is a significant decrease in oxygen tension to the right of the permeable membrane downstream. It is in this region that there is significant concentration of ATP released by RBC's uniformly across the height of channel. The second part of the present work was to determine the onset time of the release of ATP and to show complete oxygen transport profiles with changes in time. It is shown in Figure 4 that starting at initial oxygen tension of 20000 Pa (150 mmHg) at time  $t=0$  seconds, the  $PO_2$  in the channel evolves to the steady state solution at various times. The solution asymptotically approaches the steady state oxygen tension near the steady state value of 13345 Pa which is approximately equal to the 100 mmHg value (Figure 3-C) obtained in the steady state solution computed above in Maple. This is seen by the constant shaded region to the right of the permeable membrane in Figure. 8. The complete time dependent transition profiles are shown starting with  $t=5e-5$ s in Figure 5,  $t=0.02$  s in Figure 6,  $t=0.04$  s in Figure 7 and finally  $t=0.07$ s in Figure 8.

The ATP profiles did not differ between the tube and channel of radius 100 microns. For the steady-state oxygen transport and ATP concentration in an arteriole of radius 20 microns with Poiseuille flow in Figure 9, the  $PO_2$  is rather steep near the wall of the tube and in Figure 10 the concentration



of ATP is very steep near the wall of tube to a greater extent than in Fig 3-A for the tube of radius 100 microns. However the ATP levels are much lower for the realistic size of 20 microns with maximum values of 0.0009 micro-moles.(see Figure 10) Also the  $PO_2$  levels (Figure 11) far downstream are higher than those of tube of 100 microns.(see Figure 12) At the same time, the steady-state problem in an arteriole of radius 10 microns with Poiseuille flow shows as in Figure 13, that the  $PO_2$  is not as steep near the wall of tube. Also the  $PO_2$  levels (Figure 14) far downstream are lower than those of the arteriole of radius 20 microns. The results imply that there is an optimal arteriole diameter for large gradients of ATP concentration, as seen in Figure 14.

Biochemical oscillations have been studied for many years and they occur as generic processes in biological systems. The oscillation of ATP and ADP concentrations during the phosphorylation of fructose-6-phosphate(F6P), which is a key step in glycolysis is an important example of a biochemical oscillation. The enzyme phosphofructokinase (PFK) catalyzes the phosphorylation of F6P to fructose-1, 6-bisphosphate incorporating ATP as the phosphate donor[26]. To properly regulate ATP-production, ATP inhibits PFK and ADP activates PFK. If we have high ATP and low ADP then the flux of sugars into the glycol path is closed. When ATP decreases and ADP increases glycolysis recommences. Hence generated sustained oscillations are produced.

This biochemical process is associated with a rate law and accompanying kinetic parameters, where a set of rate equations describes the concentrations of all the time-varying components in the reaction network. [27].

It is noteworthy to mention that the time varying problem for ATP concentration and its temporal dynamics, as calculated in this paper is crucial in determining the ADP/ATP concentration pair which is an oscillating pair in the biochemical oscillations that occur in glycolysis in anaerobic yeast cultures.

## Conclusions

A novel approach to solve the governing equations of oxygen transport and ATP concentration in a microfluidic channel / tube with a section of permeable wall subjected to a non-zero flux condition has been presented. The results clearly show that significant concentrations of ATP are released downstream and that the onset of ATP occur at various times from zero to 70 ms and where the steady and time-dependent solutions were validated by comparison. It is shown that for smaller arterioles there is an optimal size for which greater amounts of ATP can be released. The time varying problem for ATP concentration and it's temporal dynamics, as calculated in this paper is crucial in determining the ADP/ATP concentration pair (in microchannel) which is an oscillating pair in the biochemical oscillations that occur in glycolysis in anaerobic yeast cultures. Future consideration for the present problem is to consider a network of micro-vessels and use powerful geometrical features of COMSOL Multiphysics to model and simulate ATP release in these connected micro-vessels of the microvasculature.

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## Appendix A

The following equation governs blood plasma velocity, [3]

$$v_p(y) = \frac{3 \cdot 130 \times 10^6 \frac{\mu m^3}{s}}{4 \cdot w \left( h^3 + \left( \frac{\mu_p}{\mu_c} - 1 \right) y_i^3 \right)} \begin{cases} (h^2 - y^2), & y_i \leq y \leq h \\ (h^2 - y_i^2) + \frac{\mu_p}{\mu_c} (y_i^2 - y^2), & 0 \leq y \leq y_i \end{cases}$$

where  $w$  is the width of the channel which in [6] is 1500 microns,  $y_i$  is 98 microns,  $h$  is 100 microns and  $\frac{\mu_p}{\mu_c}$  is 1.7 as obtained from [20] for a tube of radius 100 microns with hematocrit 0.2. This velocity profile was used only to compare our results to Sove's.

## Appendix B

In eq. (13) we require that the derivative,  $\partial S / \partial P$  be computed . Taking the derivative leads to the following equation, which we use in Eqs(B4,B6) and( B7) below,

$$\frac{\partial S}{\partial P} = 2.7 \cdot 3600^{2.7} \cdot \frac{P^{1.7}}{(P^{2.7} + 3600^{2.7})^2} \quad (\text{B1})$$

$$\nabla = \left[ \frac{\partial}{\partial x}, \frac{\partial}{\partial y} \right] \quad (\text{B2})$$

$$e_a \frac{\partial^2 u}{\partial t^2} + d_a \frac{\partial u}{\partial t} + \nabla \cdot \Gamma = f \quad (\text{B3})$$

$$\Gamma_x = 0 \quad (\text{B4})$$

$$\Gamma_y = -2750 \times 0.003 \cdot u_y \quad (\text{B5})$$

$$f = -1275 \times 100^2 \times 0.72 - 1147.5 \times 0.28 \times \frac{1.47}{1.33} \times \left( 1 + (0.208 / 1e - 4) \cdot \frac{\partial S}{\partial P} \right) \cdot u_x + 0.05 \cdot 2750 \times u_y \quad (\text{B6})$$

$$d_a = 6000 \cdot 0.5 \times 10^8 \cdot 0.208 \cdot \frac{\partial S}{\partial P} \quad (\text{B7})$$

## Glossary

### Nomenclature

$PO_2$  partial pressure of oxygen

$P_{50}$  partial pressure at 50% saturation

$SO_2$  oxygen saturation

$C_{ATP}$  ATP concentration

$slp$  slip velocity

$v_p$  blood plasma velocity

$v_{RBC}$  red blood cell velocity

$H_T$  tube hematocrit

$[Hb_T]$	total heme concentration
$K_p$	$O_2$ solubility in plasma
$K_{RBC}$	$O_2$ solubility in RBCs
$K_m$	$O_2$ Solubility in Porous Membrane
$D_p$	$O_2$ diffusivity in plasma
$D_{ATP}$	diffusion coefficient of ATP
$D_m$	$O_2$ Diffusivity in $O_2$ Permeable Membrane
N	Hill coefficient
R	ATP release rate
t	time t
<b>n</b>	normal unit vector
z	axial coordinate measured from channel/tube entrance
y	channel height coordinate or radial coordinate measured from tube axis
L	length of channel/tube
H	height of channel/tube

## References

- [1] Klingenberg, M., 1980, "The ADP-ATP translocation in mitochondria, a membrane potential controlled transport." *J. membrane Biol.*, **56**, pp. 97-105.
- [2] Mitchell, P., 1966, "Chemiosmotic coupling in oxidative and photosynthetic phosphorylation," *Physiol Rev.*, **41**, pp. 445-502.
- [3] Ellsworth, M.L., Ellis, C.G., Goldman, D., Stephenson, A.H., Dietrich, H.H., et al., 2009, "Erythrocytes: Oxygen sensors and modulators of vascular tone," *Physiology(Bethesda)*, **24(2)** pp. 107-118.
- [4] Dietrich, H.H., Ellsworth, M.L., Sprague, R.S., and Dacey, R.G., 2000, "Red blood cell regulation of microvascular tone through adenosine triphosphate," *Am J Physiol-Heart C*, **278(4)**, H1294.
- [5] Moschandreou, T.E., Ellis, C.G., and Goldman, D., 2011, "Influence of tissue metabolism and capillary oxygen supply on arteriolar oxygen transport: A computational model," *Math Biosci*, **232(1)**, pp. 1–10.



- [6] Sove, R.J., Ghonaim, N., Goldman, D., Ellis, C.G., 2013, "A Computational Model of a Microfluidic Device to Measure the Dynamics of Oxygen-Dependent ATP Release from Erythrocytes," *PLoS ONE*, **8(11)**, pp. 1-9.
- [7] Moschandreu, T.E., 2012, "Blood Cell- An overview of studies in hematology": Chapter 9,"RBC-ATP Theory of regulation for tissue oxygenation-ATP concentration model, ", Intech.
- [8] Erkal, J.L., Selimovic, A., Gross, B.C., Lockwood, S.Y., Walton, E.L., McNamara, S., Martin, R.S., and Spence., D.M., 2014, " 3D printed microfluidic devices with integrated versatile and reusable electrodes," *Lab on a Chip, Miniturization for chemistry, physics, biology, materials science and engineering*, Advance Article, DOI:10.1039/C4LC00171K.
- [9] Fatoyinbo, H.O., 2013," Microfluidic devices for cell manipulation,"( Edited by Li., X., and Zhou, Y.,: " Microfluidic devices for biomedical applications," pp. 283-293.
- [10] Gross, B.C., Erkal, J.L., Lockwood, S.Y., Chen C., and Spence, D.M., 2014, "Evaluation of 3D printing and it's potential impact on biotechnology and the chemical sciences," *Anal Chem*, **86(7)**, pp. 3240-3253.
- [11] Waldbaur, A., Rapp, H., Lange, K., and Rapp, B.E., 2011, "Let there be chip-towards rapid prototyping of microfluidic devices: one-step manufacturing processes," *Anal Methods*, **3(12)**, pp. 2681-2716.
- [12] Maple 13, [www.maple.ca/](http://www.maple.ca/)
- [13] COMSOL Multiphysics, [www.comsol.com](http://www.comsol.com).
- [14] Nair, P.K., Huang, N.S., Hellums, J.D., and Olson, J.S., 1990, "A simple model for prediction of oxygen transport rates by flowing blood in large capillaries," *Microvasc. Res.*, **39**, pp. 203-211.
- [15] Nair, P.K., Hellums, J.D., and Olson, J.S., 1989, "Prediction of oxygen transport rates in blood flowing in large capillaries," *Microvasc. Res.*, **38**, pp. 269-285.
- [16] Nair, P.K., 1988, "Simulation of oxygen transport in capillaries," Thesis, Rice University.
- [17] Sinha, R., 1936, *Kolloid Z.*, **76**, pp.16-24.
- [18] Pries, A.R., Secomb, T.W., and Gaehtgens, P., 1996,"Biophysical aspects of blood flow in the microvasculature," *Cardiovascular Research*, **32**, pp. 654-667.

- [19] Mayrovitz, H.N., Larnard, D., and Duda, G., 1981, "Blood velocity measurement in human conjunctival vessels," *Cardiovascular Diseases, Bulletin of the Texas Heart Institute*, **8 (4)**, pp. 509-526.
- [20] Pries, A.R., Neuhaus, D., and Gaehtgens, P., 1992, "Blood viscosity in tube flow: dependence on diameter and hematocrit," *American Journal of Physiology-Heart and Circ. Physiol.*, **263**, pp. H1770-H 1778.
- [21] Arciero, J.C., Carlson, B.E., and Secomb, T.W., 2008, "Theoretical model of metabolic blood flow regulation: roles of ATP release by red blood cells and conducted responses," *Am J. Physiol Heart C Physiol*, **295(4)**, pp. H1582-H1571.
- [22] Goldman, D., 2008, "Theoretical models of microvascular oxygen transport to tissue," *Microcirculation*, **15**, pp. 795-811.
- [23] Goldman, D., and Popel, A.S., 2001, "A computational study of the effect of vasomotion on oxygen transport from capillary networks," *J.theor. Biol.*, **209**, pp. 189-199.
- [24] Rakel, R.E., Rakel, D.P., 2011, *Textbook of family medicine*, Eighth edition, Elsevier Saunders, pp. 187
- [25] Dijkhuizen, P., Buursma, A., Fongers, T.M.E., Gerding, A.M., Oeseburg, B., and Zijlstra, W.G., 1977, "The oxygen binding capacity of human haemoglobin," *Pflugers Archiv*, **369(3)**, pp. 223-231.
- [26] McKane, A.J., Nagy, J.D., Newman, T.J., and Stefanini, M.O., 2007, "Amplified biochemical oscillations in cellular systems", *Journal of statistical physics*, **128(1/2)**, pp.165-191.
- [27] Goldbeter A., 1997, "Modelling biochemical oscillations and cellular rhythms", *Current Science* **73(11)**, pp933.